Theoretical background of the Direct Laser Interference Patterning method

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Abstract

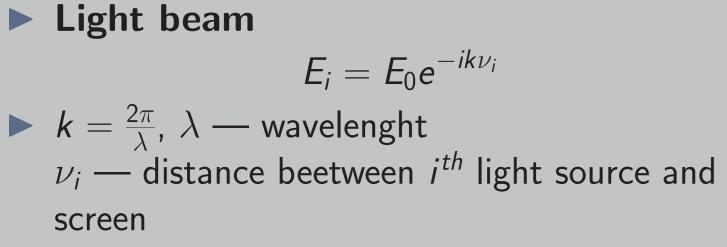
The Direct Laser Interference Patterning is a new, fast and economic method of fabrication the periodic nanostructures with size down to few nanometers. Using 2, 3 or 4 laser beams it is possible to obtain various interference patterns — from parallel periodical stripes to separated nanoislands. It can be used for obtaining new type of magnetic storage devices with separated magnetic domains.

The shape, size and period of interference image depends (apart from the number of laser beams) on optical setup geometry and used laser wavelength. On this poster there are presented theoretical calculations and simulations of interference intensity distributions for different geometrical parameters and wavelengths. There will be also shown calculation of interference image periods depended on geometry and wavelength.

I. Scheme of the experiment geometry

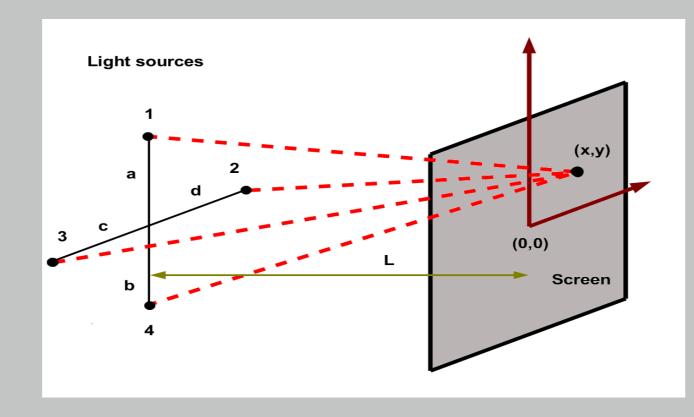
II. Two laser beams

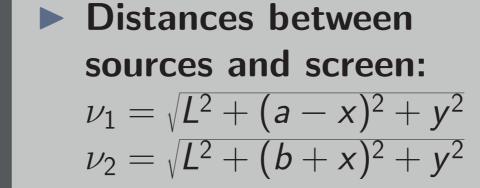
 I/I_0



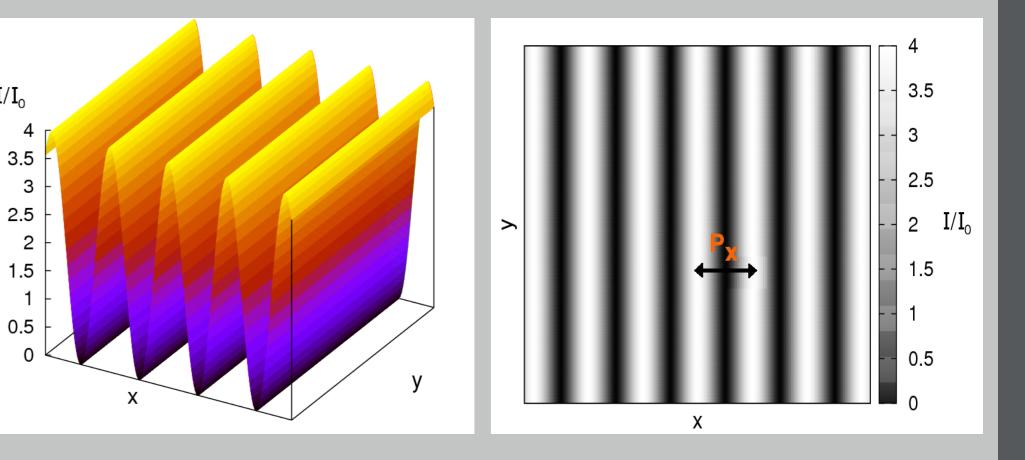
Interference intensity distribution for N beams

 $I = I_0 \left[\left(\sum_{i=1}^N \cos(k\nu_i) \right)^2 + \left(\sum_{i=1}^N \sin(k\nu_i) \right)^2 \right]$



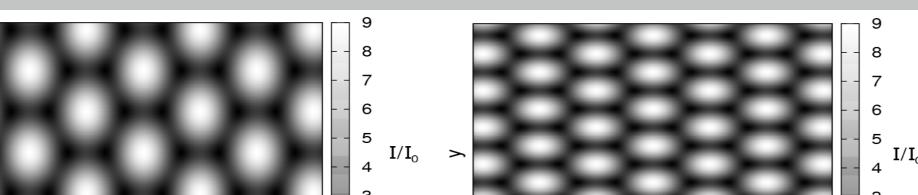


- periodical stripes,
- maximum intensity ratio $I/I_0 = 4$

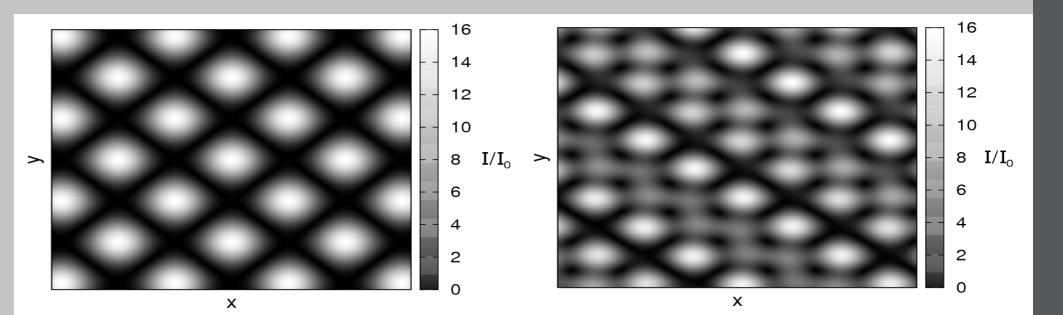


III. Three laser beams

- **Distances between** sources and screen: $\nu_1 = \sqrt{L^2 + (a - x)^2 + y^2}$ $\nu_2 = \sqrt{L^2 + (b + x)^2 + y^2}$ $\nu_3 = \sqrt{L^2 + (c - y)^2 + x^2}$
- periodical/aperiodcal pattern of dots,
- maximum intensity ratio $I/I_0 = 9$
- symmetrical configuration a = b = c
- ► asymmetrical configuration $a = b \neq c$



- **IV. Four laser beams** Distances between
 - sources and screen: $\nu_1 = \sqrt{L^2 + (a - x)^2 + y^2}$ $\nu_2 = \sqrt{L^2 + (b + x)^2 + y^2}$ $\nu_3 = \sqrt{L^2 + (c - y)^2 + x^2}$ $\nu_4 = \sqrt{L^2 + (d + y)^2 + x^2}$
- periodical/aperiodical pattern of dots,
- symmetrical configuration a = b = c = d



asymmetrical configuration $a \neq b \neq c \neq d$



maximum intensity ratio $I/I_0 = 16$

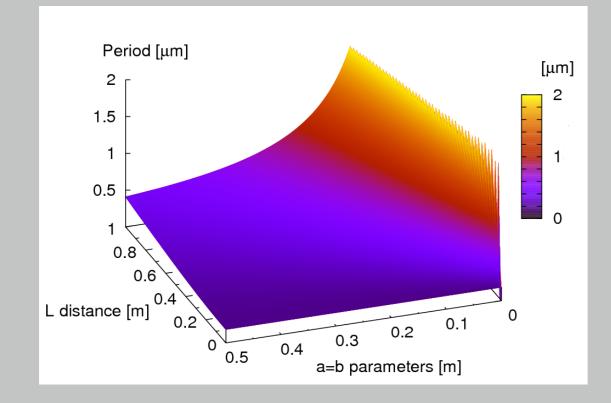
V. Period calculations for 2-beams system

Period between two stripes (n and n+1**) for symmetrical case** a = b:

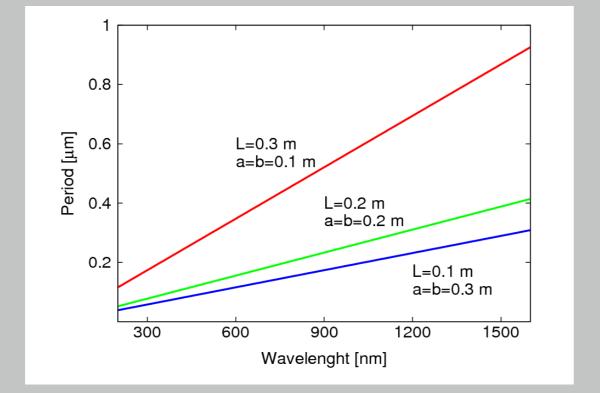
$$P(a, L, \lambda) = \left(\sqrt{\frac{4L^2\eta_{n+1} + 4a^2\eta_{n+1} - \eta_{n+1}^2}{16a^2 - 4\eta_{n+1}}} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_{n+1}^2}{16a^2 - 4\eta_n}} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_{n+1}^2}{16a^2 - 4\eta_n}} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_{n+1}^2}{16a^2 - 4\eta_n}} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_{n+1}^2}{16a^2 - 4\eta_n}} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_{n+1}^2}{16a^2 - 4\eta_n}} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_{n+1}^2}{16a^2 - 4\eta_n}} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_{n+1}^2}{16a^2 - 4\eta_n}} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_{n+1}^2}{16a^2 - 4\eta_n}} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_{n+1}^2}{16a^2 - 4\eta_n}} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_{n+1}^2}{16a^2 - 4\eta_n}} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_{n+1}^2}{16a^2 - 4\eta_n}} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_{n+1}^2}{16a^2 - 4\eta_n}} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_{n+1}^2}{16a^2 - 4\eta_n}} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_{n+1}^2}{16a^2 - 4\eta_n}} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_{n+1}^2}{16a^2 - 4\eta_n}} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_{n+1}^2}{16a^2 - 4\eta_n}} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_{n+1}^2}{16a^2 - 4\eta_n}} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_{n+1}^2}{16a^2 - 4\eta_n}} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_{n+1}^2}{16a^2 - 4\eta_n}} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_n^2}{16a^2 - 4\eta_n}} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_n^2}{16a^2 - 4\eta_n}} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_n^2}{16a^2 - 4\eta_n}} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_n^2}{16a^2 - 4\eta_n}} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_n^2}{16a^2 - 4\eta_n}} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_n^2}{16a^2 - 4\eta_n}} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_n^2}{16a^2 - 4\eta_n}} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_n^2} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_n^2}{16a^2 - 4\eta_n} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_n^2}{16a^2 - 4\eta_n} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_n^2}{16a^2 - 4\eta_n} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n + 4a^2\eta_n - \eta_n^2}{16a^2 - 4\eta_n} \right)_{n+1} - \left(\sqrt{\frac{4L^2\eta_n - \eta_n^2}{16a^2 - 4\eta_n}}$$

▶ where:

- $\eta_{n+1} = \frac{\lambda^2 (1 + 2(n+1))^2}{4} , \quad \eta_n = \frac{\lambda^2 (1 + 2n)^2}{4}$
- **Period calculated for** $\lambda = 1000 \text{ nm and}$ various *L* and *a* parameters:



Period dependent on wavelenght for fixed *L* and *a* values:



VI. Period calculations for 3-beams system

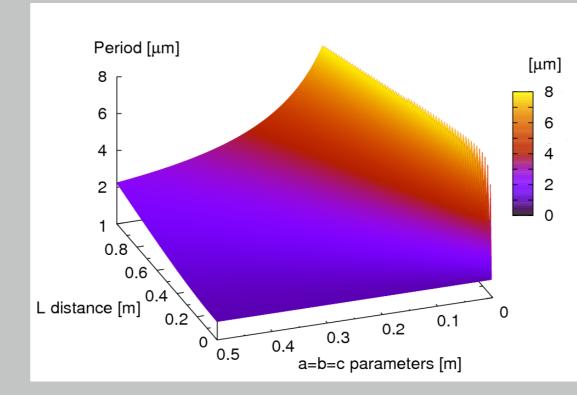
Period between two dots along the x **or** y **axis for symmetrical case** a = b = c:

$$P(a, L, \lambda) = 2 \cdot \left(\left[\sqrt{\frac{4\phi_{n+1}^2 L^2 + 4\phi_{n+1}^2 a^2 - \phi_{n+1}^4}{16a^2 - 4\phi_{n+1}^2}} \right]_{n+1} - \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^2 a^2 - \phi_n^4}{16a^2 - 4\phi_n^2}} \right]_{n+1} \right]_{n+1} - \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^2 a^2 - \phi_n^4}{16a^2 - 4\phi_n^2}} \right]_{n+1} + \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^2 a^2 - \phi_n^4}{16a^2 - 4\phi_n^2}} \right]_{n+1} + \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^2 a^2 - \phi_n^4}{16a^2 - 4\phi_n^2}} \right]_{n+1} + \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^2 a^2 - \phi_n^4}{16a^2 - 4\phi_n^2}} \right]_{n+1} + \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^2 a^2 - \phi_n^4}{16a^2 - 4\phi_n^2}} \right]_{n+1} + \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^2 a^2 - \phi_n^4}{16a^2 - 4\phi_n^2}} \right]_{n+1} + \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^2 a^2 - \phi_n^4}{16a^2 - 4\phi_n^2}} \right]_{n+1} + \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^2 a^2 - \phi_n^4}{16a^2 - 4\phi_n^2}} \right]_{n+1} + \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^2 a^2 - \phi_n^4}{16a^2 - 4\phi_n^2}} \right]_{n+1} + \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^2 a^2 - \phi_n^4}{16a^2 - 4\phi_n^2}} \right]_{n+1} + \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^2 a^2 - \phi_n^4}{16a^2 - 4\phi_n^2}} \right]_{n+1} + \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^2 a^2 - \phi_n^4}{16a^2 - 4\phi_n^2}} \right]_{n+1} + \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^2 a^2 - \phi_n^4}{16a^2 - 4\phi_n^2}} \right]_{n+1} + \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^2 a^2 - \phi_n^4}{16a^2 - 4\phi_n^2}} \right]_{n+1} + \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^2 a^2 - \phi_n^4}{16a^2 - 4\phi_n^2}} \right]_{n+1} + \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^2 a^2 - \phi_n^4}{16a^2 - 4\phi_n^2}} \right]_{n+1} + \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^2 a^2 - \phi_n^4}{16a^2 - 4\phi_n^2}} \right]_{n+1} + \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^2 a^2 - \phi_n^4}{16a^2 - 4\phi_n^2}} \right]_{n+1} + \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^2 a^2 - \phi_n^4}{16a^2 - 4\phi_n^2}} \right]_{n+1} + \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^2 a^2 - \phi_n^4}{16a^2 - 4\phi_n^2}} \right]_{n+1} + \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^2 a^2 - \phi_n^4}{16a^2 - 4\phi_n^2}} \right]_{n+1} + \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^2 a^2 - \phi_n^4}{16a^2 - 4\phi_n^4}} \right]_{n+1} + \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^2 a^2 - \phi_n^4}{16a^2 - 4\phi_n^4}} \right]_{n+1} + \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^2 a^2 - \phi_n^4}{16a^2 - 4\phi_n^4} \right]_{n+1} + \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^4 - \phi_n^4}{16a^2 - 4\phi_n^4} \right]_{n+1} + \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^4 - \phi_n^4}{16a^2 - 4\phi_n^4} \right]_{n+1} + \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^4 - \phi_n^4}{16a^2 - 4\phi_n^4} \right]_{n+1} + \left[\sqrt{\frac{4\phi_n^2 L^2 + 4\phi_n^4 - \phi_n^4 - \phi_n^4}$$

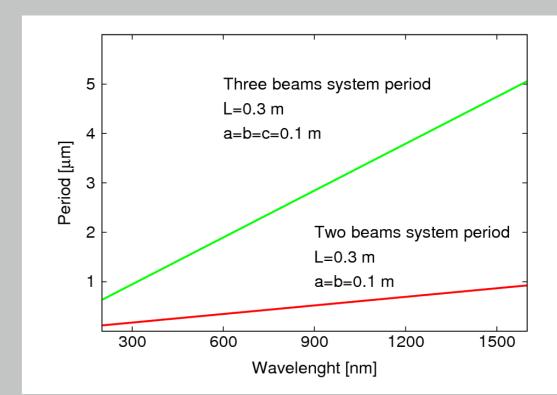
▶ where:

$$\phi_{n+1} = (n+1)\lambda$$
, $\phi_n = n\lambda$

Period calculated for $\lambda = 1000 \text{ nm}$ and various *L* and *a* parameters:



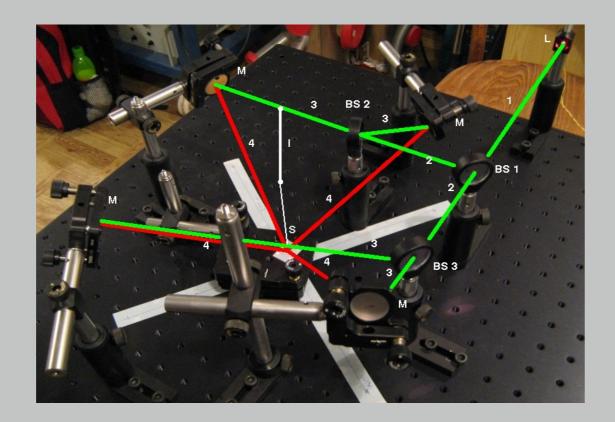
Comparison the periods for 2- and 3-beams systems:



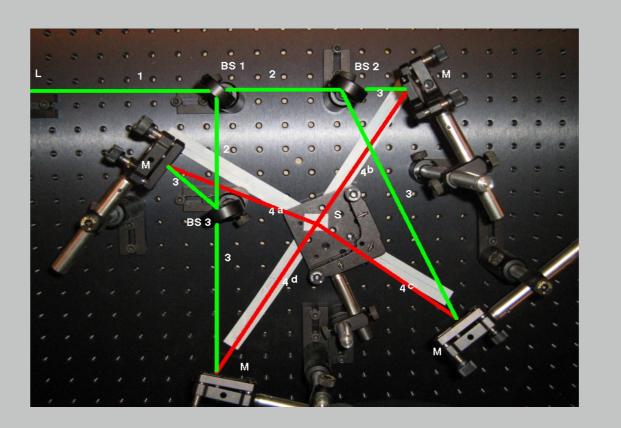
VII. The experimental setup and example of calculated interference pattern

Used optical elements:

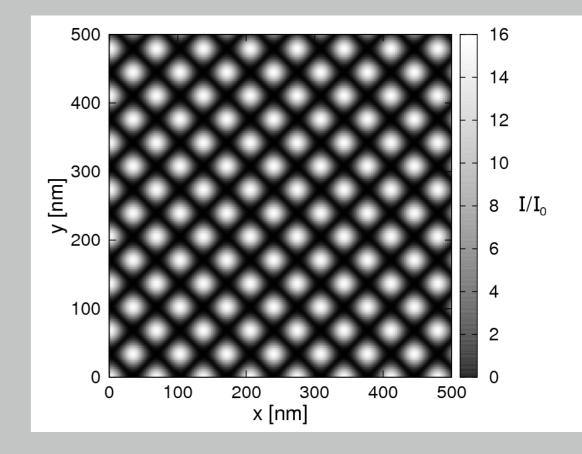
- ▶ four mirrors (M1 M4),
- \blacktriangleright three 50% 50% beam splitters (BS1 - BS3),
- holders for optical elements,
- laser with wavelenght 600 nm (L)
- **Interference geometry:**
- ► four laser beams,
- ▶ symmetrical case a = b = c = d = 0,31 m, the L-distance between sample (screen) and
- plane of light sources was L = 0, 17 m
- **Experimental setup** side view



Experimental setup — top view



Calculated intereference pattern



http://www.ifj.edu.pl, http://www.ifj.edu.pl/dept/no5/nz53/

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